Modeling of particulate matter transport in atmospheric boundary layer following dust emission from source areas

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ABSTRACT

A two-dimensional model for particulate matter (PM) dispersion due to dust emission from soils is presented. Field experiments were performed at a dust source site (Negev loess soil) with a portable boundary layer wind tunnel to determine the emitted PM fluxes for different wind speeds and varying soil conditions. The numerical model is formulated using parameterizations based on the aeolian experiments. The wind velocity profiles used in the simulations were fitted from data obtained in field measurements. Size distribution of the emitted dust particles in the numerical simulations was taken into account using a Monte Carlo method. The PM concentration distributions at a distance of several kilometers from the dust source under specific shear velocities and PM fluxes from the soil were determined numerically by solving advection–diffusion equation. The obtained PM\(_{10}\) concentrations under typical wind and soil conditions are supported by PM data recorded over time in a standard environmental monitoring station. The model enhances our capacity of quantification of dust processes to support climate models as well as health risk assessment.

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1. Introduction

Mineral dust is a key agent involved in a wide range of physical, chemical and biological processes of the Earth system. Dust particles profoundly affect the energy, carbon and water cycles of the climate system (see e.g., Shao et al., 2011). It has been recognized that dust aerosol can also have significant impacts on human health. During dust storms, concentrations of dust particles having aerodynamic diameter less or equal to 10 \(\mu\)m (PM\(_{10}\)) in arid areas can exceed significantly the World Health Organization (WHO) guideline for air quality (Katra et al., 2014a), reaching outdoor and indoor concentrations in arid environments as high as 5000 and 1500 \(\mu\)g/m\(^3\), respectively (Krasnov et al., 2015). Other studies highlighted the role of PM from natural dust as an important environmental pollutant for human health impacts (Vodonos et al., 2014; Yitshak-Sade et al., 2014).

Modeling of dust transport is of profound importance for understanding the dust cycle at different scales. It requires combined theoretical and experimental study of particle emission (in consequence of wind soil erosion) and dust transport in the atmospheric boundary layer (ABL). In particular, modeling of atmospheric dust dispersion requires reliable data on measured input variables, such as grain size distribution, wind speed, sediment properties, dust-emission rate (see, e.g., Gillies et al., 2006; Shao, 2008; Durán et al., 2011; Kok et al., 2012).

Understanding the role of dust in the Earth system has prompted intensive development of dust emission models since the late 1980s (Shao et al., 2011). These developments spurred the intensive efforts in the modeling of dust transport. Pilinis et al. (1987) developed a mathematical model that describes the evolution of size and chemical composition distribution of atmospheric aerosols based on a sectional representation of the size distribution, and treats dynamics and thermodynamics of multicomponent atmospheric aerosols. Lu and Shao (1999) developed a theoretical model for the prediction of dust emission rate caused by saltation bombardment, based on the dust volume removal caused by impacting sand grains as they plough into the soil surface. In the present study, the dust emission generated by saltation bombardment has been modeled from the perspective of volume removal by saltating particles as proposed by Lu and Shao (1999). As Lu and Shao (1999), Shao (2001) noted, more field measurements are required in order to verify this approach.
In spite of intensive development of advanced dust emission models, there is a gap in quantification of dust transport following dust emission from soils. Moreover, the knowledge about the impact of surface-property variability on dust fluxes from source areas is still severely lacking (Katra and Lancaster, 2008).

For individual wind-erosion events, wind shear near the surface is responsible for particle entrainment into the atmosphere, and turbulence in the atmospheric boundary layer is important for particle diffusion and deposition (Shao, 2008). Many of the early experiments on wind erosion (see e.g., Bagnold, 1936; Kawamura, 1951; Lettau, 1969; Lettau and Lettau, 1977; Nickling and Gillies, 1993; Bottema, 1996; MacDonald et al., 1998) show that in the case of negligible wake interference between the surface obstacles the mean velocity profile approaching each obstacle is logarithmic.

In this study, aeolian erosion under different wind and soil conditions in an area located in the northern Negev Desert (Israel) was simulated for modeling dust PM dispersion in the atmospheric boundary layer at a distance of several kilometers from the dust source. Field experiments were used to determine a set of parameters for the dust dispersion model that include dust emission fluxes and velocity profiles.

2. Mean wind velocity profile

The mean wind velocity profile for the studied case is required for simulating the PM dispersion after emission of dust from the soil. To this end we use standard equations of atmospheric boundary layer theory (see e.g., Shao, 2008). Goossens and Offer (1990) showed that in the Northern Negev area the depth of the Atmospheric Boundary Layer (ABL) is on the order of 500–600 m. Since we consider horizontal transport of dust particles over large distances, the assumptions of the validity of the boundary layer approximation are satisfied. Offer and Goossens (1994) showed that in the lowermost 15% of the ABL, the wind profile (in neutral atmospheric conditions) can be described either by the semi-empirical logarithmic law (see, e.g., Oke, 1987):

$$u = \frac{u_1}{\kappa} \ln \left( \frac{z}{z_0} \right)$$

(1)

or by the power law (see, e.g., Offer and Goossens, 1994):

$$u_1 = u_2 \left( \frac{z_1}{z_2} \right)^{2/3}$$

(2)

In Eq. (1) $u_1 = \sqrt{\tau / \rho_s} \sim \sigma$ is friction velocity, which is a scaling parameter proportional to the velocity gradient in boundary layer flow; $\tau$ is the shear stress at the surface level and $\rho_s$ is air density; $\sigma$ is standard deviation of velocity fluctuations (see e.g., Bagnold, 1941; Shao, 2008; Kok et al., 2012); $\kappa$ is the von Karman constant, $\kappa = 0.35 – 0.4$; $z_0$ is the aerodynamic surface roughness length.

<table>
<thead>
<tr>
<th>Height, $h$ (m)</th>
<th>Average wind velocity, $u$ (m/s)</th>
</tr>
</thead>
<tbody>
<tr>
<td>0.68</td>
<td>4.15</td>
</tr>
<tr>
<td>1.18</td>
<td>4.49</td>
</tr>
<tr>
<td>2.0</td>
<td>4.76</td>
</tr>
<tr>
<td>3.36</td>
<td>5.04</td>
</tr>
<tr>
<td>5.64</td>
<td>5.31</td>
</tr>
<tr>
<td>9.43</td>
<td>6.09</td>
</tr>
</tbody>
</table>

Table 1
Mean wind velocity vs. height based on field experiment of 12 June 2014.
that is approximately equal 1/30 of the field roughness elements height in turbulent flow. In the Negev area roughness elements’ height (pebbles and shrubs) is of the order 0.2–0.3 m (Offer and Goossens, 1994). In Eq. (2) \( u_1 \) and \( u_2 \) are the mean wind speeds at heights \( z_1 \) and \( z_2 \), respectively. For the Negev area the value of the power law exponent \( \alpha \) is in the range 0.16–0.17 (Offer and Goossens, 1994).

For this study, measurements of wind profiles were conducted in the Northern Negev with a wind mast in conjunction with experiments to quantify dust emission from the soils (see more details in the next section). A 10-meter wind mast was equipped with 6 cup anemometers positioned at: 0.68, 1.18, 2.0, 3.36, 5.64 and 9.43 m above ground level (AGL). The instruments record wind speeds in the range of 0–50 m/s with the accuracy of ±0.49 m/s and with a time interval of 1 min. The anemometers were calibrated in the laboratory before the field measurements. Wind speed measurements are shown in Fig. 1.

For each height the average wind speed was calculated as follows:

\[
\bar{u} = \frac{1}{T - t_0} \int_{t_0}^{T} u(t) \, dt, \tag{3}
\]

where \( t_0 \) is the initial time and \([t_0, T]\) is the time interval during which measurements were conducted. Average wind velocity at each height was determined by numerical integration of wind

**Table 2**

Results of particle fluxes from soils.

<table>
<thead>
<tr>
<th>Area</th>
<th>Open area</th>
<th>Disaggregated</th>
<th>Nature preserve</th>
<th>Aggregated</th>
<th>Disaggregated</th>
</tr>
</thead>
<tbody>
<tr>
<td>Condition (soil)</td>
<td>Aggregated</td>
<td>Disaggregated</td>
<td>Aggregated</td>
<td>Disaggregated</td>
<td></td>
</tr>
<tr>
<td>Frequency (Hz)</td>
<td>29</td>
<td>44</td>
<td>29</td>
<td>44</td>
<td>29</td>
</tr>
<tr>
<td>PM(_{2.5}) (mg m(^{-2}) min(^{-1}))</td>
<td>17.1</td>
<td>256.5</td>
<td>52.3</td>
<td>895.3</td>
<td>1.2</td>
</tr>
<tr>
<td></td>
<td>120.9</td>
<td>8.3</td>
<td>120.9</td>
<td>8.3</td>
<td>120.9</td>
</tr>
</tbody>
</table>
speed time series $u(t)$ using trapezoidal integration rule, and the obtained results are presented in Table 1.

The above dependence of the mean wind velocity vs. height was approximated by a logarithmic profile (see Eq. (1)) using the least squares method. Therewith the best fit to the observations is determined for the friction (shear) velocity $u_\tau = 0.27 \text{ m/s}$ and the roughness $z_0 = 1.6 \times 10^{-3}$ m with 95% confidence bounds. The results of calculations of the mean wind velocity as function of altitude are showed in Fig. 2.

In Fig. 2 the mean wind speed profiles were calculated using two different values of roughness, $z_0 = 4.0 \mu m$ and $z_0 = 10.0 \mu m$. Roughness $z_0 = 4.0 \mu m$ is calculated as 1/30 of the roughness elements corresponding to the large particles with the size of 120 $\mu m$ (Bagnold, 1941), while the roughness $z_0 = 10.0 \mu m$ corresponds to the roughness elements such as rocks and shrubs at the Negev area (see Offer and Goossens, 1994). The values of shear velocity for these aerodynamic roughness values are $u_\tau = 0.15 \text{ m/s}$ and $u_\tau = 0.36 \text{ m/s}$, respectively. The double dot-dashed line in Fig. 2 describes the power law (3) where $u_2 = 4.76 \text{ m/s}$ is the magnitude of the mean wind velocity at the height $z_2 = 2 \text{ m}$ (see Table 1).

3. Aeolian experiments on dust emission

Boundary-layer wind tunnels enable aeolian simulations under standardized quasi-natural wind conditions (Leys and Raupach, 1991; Shao, 2008) and provide quantitative information on aeolian particle transport including sand fluxes (Katra et al., 2014b) and dust emission rates from soils (Houser and Nickling, 2001; Sharratt et al., 2010; Singh et al., 2012; Van Pelt et al., 2013). The wind-tunnel experiments were conducted in the loess soil of the northern Negev. Many of the loess soils throughout the world are associated with wind erosion activity, which is largely a result of environmental changes such as droughts and increased human activities, thereby amplifying dust emissions from the disturbed surfaces. These soils contain a significant fraction of silt and clay (<50 $\mu m$), 71.6% in nature preserve and 64.2% in open area. The presence of a sand fraction in its loamy texture enables the entrainment of fine particles by ballistic impact (saltation bombardment) (Kok et al., 2012). The measurements were conducted in experimental plots of bare and dry soils representing common surfaces in the Northern Negev: open area – an area with human activities that is mainly grazing, and a nature preserve – where human activities are highly restricted. In both cases the soils were tested in two different conditions of the topsoil: undisturbed (aggregated), and disturbed (disaggregated) conditions. In the latter condition, the soils were treated in the field to simulate a short-term disturbance of human activities by mechanical operation (Bacon et al., 2011) for a reduced soil aggregation. The portable boundary-layer wind tunnel has been designed to study dust emission in the field by simulating aeolian processes (Fig. 3). The wind tunnel has a cross sectional area of 0.5 x 0.5 m², with an open-floored working section up to 10 m long. Air push or air suction flow in the tunnel is generated by an axial fan up to maximum speed of 18 m s⁻¹ measured along the centerline. Instruments installed in the test section of the tunnel enable quantification of: (1) the vertical wind speed profile for calculating friction velocity and aerodynamic roughness height, (2) salination flux, and (3) dust concentration profile including PM₁₀.

The wind tunnel was operated at the experimental plots at fan frequencies (wind speeds) of 23 Hz (4 ms⁻¹) and 44 Hz (9 ms⁻¹), which represent the minimum wind speed required for aeolian erosion and an average wind speed in dust storms in the study area, respectively. For each fan frequency the wind speed was measured during testing for the specific soils. The measured vertical velocity profiles in the wind tunnel (see Fig. 4) exhibit the same semi-log form as velocity profiles determined in the experiments with a wind mast (see Table 1 and Fig. 2).

At the beginning of each test, the concentration of the suspended particles in the wind tunnel was measured by a real-time dust monitor for background PM₁₀ values (DustTrak 8534, TSI). The DustTrak installed in the test section (Fig. 3) recorded PM concentration ($\mu g \text{ m}^{-3}$) at 1 Hz. Each experiment (surface/condition/frequency) was performed in three replicates with a run time of 600 s. The recorded PM₁₀ data (concentration) was converted into mass flux emitted from the soil surface ($F_{PM}$) based on the wind tunnel dimensions and area of the experimental plot:

$$F_{PM} = C_{PM} V_t / (A_p t)$$

where $C_{PM}$ is the recorded PM concentrations (mg m⁻³), $V_t$ is the volume air in the wind tunnel (1.85 m³), $A_p$ is the area of the experimental plot (3.75 m²), and t is time (in minutes).

The obtained results (see Table 2) demonstrate the dependence of dust emission in surface variability and wind speed. Overall, there were higher dust fluxes from the open area as compared to the natural preserve. The open area is subjected to grazing activities, which impact the topsoil properties and reduce the surface stability against erosion. Analyses of the topsoil properties show lower contents of soil cementing material (e.g., soil organic matter, clays, calcium carbonate) in the open area compared to the natural preserve and therefore lower content of large aggregates. Such differences in the soil properties between grazing and natural plots have been demonstrated in previous studies (see e.g., Singh et al., 2012; Tanner et al., 2016). It is generally assumed that soils with a greater amount of large aggregates have stronger resistance against erosion. Only a few studies referred to the soil aggregation in aeolian processes, but with a focus on the wind erodible fraction (EF) (aggregates in diameter of <840 $\mu m$) (Zobeck et al., 2013). In both experimental plots, higher wind speed (corresponding to larger fan frequency) and disturbed topsoil (disaggregation) increase the PM₁₀ fluxes from the soil. These results were used for parameterization of dust fluxes in the numerical simulations.

4. Numerical simulations

Atmospheric dust concentration at given time and location can be described by the advection–diffusion equation for dust particles (see e.g., Shao et al., 2011):
depends on the vertical coordinate only (see Eq. (1)). In the steady-state case the Eq. (5) can be rewritten as follows:

\[ u \frac{\partial c}{\partial x} - w_t \frac{\partial c}{\partial z} = \frac{\partial}{\partial z} \left( K_{pz} \frac{\partial c}{\partial z} \right) \] (6)

The boundary conditions for Eq. (6) read:

\[ w_t c - K_{pz} \frac{\partial c}{\partial z} = Q_0 \left[ 1 - H(x - x_5) \right] \text{ at } z = 0 \] (7)

\[ K_{pz} \frac{\partial c}{\partial z} = 0 \text{ at } z = h \] (8)

\[ c(0, z) = 0. \] (9)

where \( h \) is the height of the ABL. In the calculations we assumed that \( h = 500 \text{ m} \).

In Eqs. (6)–(9) \( Q_0 \) is the net dust flux at the surface, \( H(x) \) is Heaviside step function, \( x_5 \) is the length of the source area and \( w_t \) is the terminal settling velocity of deposited dust particles (Seinfeld and Pandis, 2006, p. 408):

\[ w_t = \frac{1}{18} \frac{d_p^2 \rho_p g C_c}{\mu}, \] (10)

where \( d_p \) is the diameter of the dust particle, \( \rho_p \) is the density of the dust particle; \( g \) is gravitational acceleration; \( C_c \) is the Cunningham slip correction factor; \( \mu \) is the coefficient of the dynamic viscosity of the air. Estimations show that for dust particles with the diameter \( d_p = 10 \mu \text{m} \) and density \( \rho_p = 2.6 \text{ g/cm}^3 \) the terminal settling in air velocity is \( \approx 7.9 \cdot 10^{-3} \text{ m/s} \). For a dust particle with the diameter \( d_p = 120 \mu \text{m} \) the terminal velocity is on the order of \( 1 \text{ m/s} \).

The coefficient of eddy diffusivity can be calculated as follows (see e.g., Shao, 2008):

\[ K_{pz} = \kappa u z. \] (11)
Fig. 7. Distribution of dust particles concentration calculated in XZ-plane at $Y = 0$ (friction velocity $u_* = 0.27$ m/s, aerodynamic roughness $z_0 = 1.6 \times 10^{-3}$ m, dust flux at the surface $Q_0 = 17.1$ mg m$^{-2}$ min$^{-1}$ – non-cohesive state): (a) diameter of the particles $d_p = 10$ $\mu$m; (b) log-normal distribution of particle diameter is taken into account ($d_p = 10$ $\mu$m, $\sigma_d/d_p = 0.3$).

Fig. 8. Distribution of dust particles concentration calculated in XZ-plane at $Y = 0$ (friction velocity $u_* = 0.27$ m/s, aerodynamic roughness $z_0 = 1.6 \times 10^{-3}$ m, log-normal distribution of particle diameter is taken into account – $d_p = 5$ $\mu$m, $\sigma_d/d_p = 0.3$); (a) dust flux at the surface $Q_0 = 52.3$ mg m$^{-2}$ min$^{-1}$ (non-cohesive state); (b) dust flux at the surface $Q_0 = 17.1$ mg m$^{-2}$ min$^{-1}$ (cohesive state).
Transport of dust particles by turbulent diffusion as well as dust particles deposition strongly depends on the dust particle diameter. In our calculations we assumed a log-normal size distribution for the dust particles:

\[
f_{d_p} = \frac{1}{\sigma \sqrt{2\pi}} e^{\frac{-(\ln d_p - \mu)^2}{2\sigma^2}}
\]

Parameters \(\mu\) and \(\sigma\) of the log-normal distribution are determined by \(d_r\) and \(\sigma_r\), mean and variance of particle size distribution, respectfully:

\[
\sigma = \ln \left( \frac{\sigma_r^2 + 1}{d_r^2} \right), \quad \mu = \ln d_r - \frac{\sigma^2}{2}
\]

These parameters can be estimated as follows:

\[
d_r = \frac{1}{N} \sum_{i=1}^{N} d_{p,i}, \quad \sigma_r^2 = \frac{1}{N-1} \sum_{i=1}^{N} (d_{p,i} - d_r)^2,
\]

where \(d_{p,i}\) are measured diameters of dust particles and \(N\) is the number of measurements.

Numerical analysis of a dust transport model described by the advection–diffusion equation (6) supplemented with the initial and boundary conditions (7)–(9) was performed using MATLAB numerical toolbox. Particle size distribution was taken into account using the Monte Carlo method whereby the particle concentration distribution was calculated by solving the initial boundary-value problem given by Eq. (12). The particle concentration distribution was determined by averaging over an ensemble of 1000 particle concentration distributions determined for randomly chosen particle diameters. The Monte Carlo algorithm was implemented into a MATLAB routine and integrated into the dust transport model.

5. Results and discussion

The model was applied to the calculations of dust emission and transport from different area sources with lengths 3, 5 and 10 km. Calculations were performed for particles with average
aerodynamic diameters 5, 10 and 20 μm as well as for PM₅, PM₁₀ and PM₂₀. In numerical simulations we used Monte-Carlo method whereby the advection–diffusion equation was solved for 1000 particle diameters sampled from the log-normal distribution given by Eq. (12). The particle concentration distribution was determined as average of these solutions. It was assumed that the wind velocity profile is determined by Eq. (1) where the friction velocity and aerodynamic roughness are determined by fitting the experimental data \( u_0 = 0.27 \text{ m/s} \) and \( z_0 = 1.6 \times 10^{-3} \text{ m} \) and the measured dust fluxes \( Q₀ = 52.3 \text{ mg m}^{-2} \text{ min}^{-1} \) and \( Q₀ = 17.1 \text{ mg m}^{-2} \text{ min}^{-1} \) in open area for the aggregated and disaggregated conditions, respectively (see Table 2).

Results of calculation of dust (PM₁₀) concentration profiles are shown in Fig. 5. In this case the length of the source area was assumed to be 10 km. Calculations were performed for dust flux \( Q₀ = 52.3 \text{ mg m}^{-2} \text{ min}^{-1} \), and different distances from the origin of the source area, namely 0.1, 0.5, 1, 2, 3, 4, 5 and 10 km. Inspection of the obtained vertical concentration profiles shows that the maximum concentration of PM₁₀ at the ground exceeds 35 mg m⁻³.

Results of calculations of the concentration distribution of the dust particles \( dp = 10 \mu \text{m} \) near the ground for two values of dust flux \( Q₀ = 52.3 \text{ mg m}^{-2} \text{ min}^{-1} \) and \( Q₀ = 17.1 \text{ mg m}^{-2} \text{ min}^{-1} \) are shown in Figs. 6 and 7, respectively.

Calculations for monosize particles with the diameter 10 μm and dust fluxes \( Q₀ = 52.3 \text{ mg m}^{-2} \text{ min}^{-1} \) and \( Q₀ = 17.1 \text{ mg m}^{-2} \text{ min}^{-1} \) are presented in Figs. 6a and 7a. Figs. 6b and 7b present the results of calculations taking into account the log-normal distribution of particle diameter (in this case \( d₀ = 10 \mu \text{m} \), \( σ_r/d₀ = 0.3 \)) and for dust fluxes \( Q₀ = 52.3 \text{ mg m}^{-2} \text{ min}^{-1} \) and \( Q₀ = 17.1 \text{ mg m}^{-2} \text{ min}^{-1} \), respectively.

Dust particle concentrations decrease with height and increase in \( x \)-direction at low height level (near the ground) along the linear dust source (see Figs. 6 and 7). Comparison of the results of calculations (see Fig. 6a and b and Fig. 7a and b) shows that neglecting of particle size distribution may result in underestimation of particles concentrations in vertical direction.

Results of calculations of the concentration distribution of the dust particles taking into account particles size distributions for different values of the mean are shown in Figs. 8 and 9. Calculations were performed for the two values of dust flux \( Q₀ = 52.3 \text{ mg m}^{-2} \text{ min}^{-1} \) and \( Q₀ = 17.1 \text{ mg m}^{-2} \text{ min}^{-1} \). Parameters \( d₀ \) and \( σ_r \) of log-normal distribution (see Eq. (12)) were assumed...
to be 5 μm and 20 μm and $\sigma = 0.3 d$, correspondingly. The results revealed that the concentration of the dust particles strongly depends on the size of the particles.

Results of calculations of dust (PM$_{10}$) concentration profiles for a source area $x_S = 5$ km are shown in Fig. 10. The dust flux was assumed to be $Q_0 = 52.3$ mg m$^{-2}$ min$^{-1}$, and we used a logarithmic wind velocity profile corresponding to friction velocity $u_* = 0.27$ m/s and aerodynamic roughness $z_0 = 1.6 \times 10^{-3}$ m.

Results of calculations of the concentration distribution of PM$_{10}$ near the ground emitted from source area with the length 5 km are shown in Fig. 11. Note that particles are advected beyond the length of the source area, and PM concentration distribution sharply changes outside the source area.

We also performed calculation of dust (PM$_{10}$) concentration profiles and concentration distribution of PM$_{10}$ near the ground emitted from source area with length of $x_S = 3$ km. The results are shown in Figs. 12 and 13, correspondingly.

The study clearly demonstrates the impact of surface-property variability on dust fluxes (Table 2) and thus on the results of dust and PM transport models. The performed numerical calculations (see Figs. 5–13) show that at the distance of several kilometers from the dust source there are high atmospheric particulate matter (PM) concentrations in terms of air quality. According to the WHO guidelines the low end of the range of concentrations at which adverse health effects has been demonstrated is not greatly above the background concentration, which for particles smaller than 10 μm (PM$_{10}$) has been estimated to be 0.050 mg m$^{-3}$. Note that the high concentrations due to dust emission from this source area (bare and dry loess soil in a semi-arid zone) occurs even under the conditions of relatively low friction velocity (0.27 m/s) and undisturbed soils (dust flux at the surface $Q_0 = 17.1$ mg m$^{-2}$ min$^{-1}$). At 10 m height, the results show about 2.5 mg m$^{-3}$ and 4.5 mg m$^{-3}$ for PM$_{10}$ and PM$_{2.5}$, respectively (Figs. 6b and 7b). The simulation results of PM$_{10}$ concentrations were compared with the recorded data in a monitoring station located at a distance of about 10 km west from the dust source (open area, Table 2). The monitoring station records PM$_{10}$ data at height of 10 m by a dichotomous ambient particulate monitor (Thermo Scientific 1405DF; Thermo Fisher Scientific Inc.), which provides a continuous direct measurement of particulate monitor (TEOMs). PM$_{10}$ analysis for the period of 2001–2012 enabled assessment of dust concentration distribution. Moreover, such models can improve our understanding on loss of soil resources such as in agricultural fields due to wind erosion as well as on air pollution and exposure risks in urban environments located close to dust sources.

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